

NonRelativistic Quantum Mechanics Formulas

$$|x_{\pm}\rangle = \frac{|+\rangle \pm |-\rangle}{\sqrt{2}} \quad |y_{\pm}\rangle = \frac{|+\rangle \pm i|-\rangle}{\sqrt{2}} \quad \text{Norm: } \sqrt{\langle\alpha|\alpha\rangle} \quad \text{Completeness: } \sum_i |i\rangle\langle i| = 1$$

$$\vec{S} = \frac{\hbar}{2}\vec{\sigma} \quad \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\langle(\Delta A)^2\rangle\langle(\Delta B)^2\rangle \geq \frac{1}{4}|\langle[A, B]\rangle|^2$$

Change of Basis: If $|b^{(i)}\rangle = U|a^{(i)}\rangle$ where $U = \sum_i |b^{(i)}\rangle\langle a^{(i)}|$ then

Column vectors: (New) = (U^\dagger) (Old) Operators: $X' = U^\dagger X U$

$$[x_i, x_j] = 0, \quad [x_i, p_j] = i\hbar\delta_{ij} \quad [p_i, p_j] = 0 \quad \langle\vec{x}|\vec{p}\rangle = \frac{1}{(2\pi\hbar)^{3/2}} \exp\left(\frac{i\vec{p}\cdot\vec{x}}{\hbar}\right)$$

$$U(t, t_0) = T \left[\exp\left(-\frac{i}{\hbar} \int_{t_0}^t dt' H(t')\right) \right]$$

$$A^{(H)}(t) = \mathcal{U}^\dagger(t) A^{(S)} \mathcal{U}(t) \quad \frac{dA^{(H)}}{dt} = \frac{i}{\hbar} [H, A^{(H)}]$$

$$a = \sqrt{\frac{m\omega}{2\hbar}} \left(\hat{x} + \frac{i\hat{p}}{m\omega} \right) \quad [a, a^\dagger] = 1 \quad \int_{x_1}^{x_2} dx \sqrt{2m(E - V(x))} = \left(n + \frac{1}{2}\right)\pi\hbar$$

NonRelativistic Quantum Mechanics Formulas, contd.

$$\int_{-1}^1 P_m(x)P_n(x)dx = \frac{2\delta_{mn}}{2n+1}$$

$$[J_z, T_q^{(k)}] = \hbar q T_q^{(k)} \qquad [J_{\pm}, T_q^{(k)}] = \hbar \sqrt{k(k+1) - q(q \pm 1)} T_{q \pm 1}^{(k)}$$

$$\mathcal{D}^\dagger(R)T_q^{(k)}\mathcal{D}(R) = \sum_{q'=-k}^k \mathcal{D}_{qq'}^{(k)*}(R)T_{q'}^{(k)} \qquad [\vec{J} \cdot \hat{n}, T_q^{(k)}] = \sum_{q'} T_{q'}^{(k)} \langle k q' | \vec{J} \cdot \hat{n} | k q \rangle$$

$$T_q^{(k)} = \sum_{q_1} \sum_{q_2} \langle k_1 k_2; q_1 q_2 | k q \rangle X_{q_1}^{(k_1)} Z_{q_2}^{(k_2)}$$

$$\langle \alpha', j' m' | T_q^{(k)} | \alpha, j m \rangle = \langle j k; m q | j k; j' m' \rangle \frac{\langle \alpha' j' || T^{(k)} || \alpha j \rangle}{\sqrt{2j+1}}$$

$$a = \sqrt{\frac{m\omega}{2\hbar}} \left(x + \frac{i p}{m\omega} \right) \qquad x = \sqrt{\frac{\hbar}{2m\omega}} (a + a^\dagger) \qquad [a, a^\dagger] = 1$$

$$H = \hbar \omega \left(a^\dagger a + \frac{1}{2} \right) \qquad a | n \rangle = \sqrt{n} | n-1 \rangle \qquad a^\dagger | n \rangle = \sqrt{n+1} | n+1 \rangle$$

$$\Pi^\dagger \vec{x} \Pi = -\vec{x} \qquad \Pi | \alpha, \ell m \rangle = (-1)^\ell | \alpha, \ell m \rangle \qquad \Theta | \ell, m \rangle = (-1)^m | \ell, m \rangle$$

$$\Theta \vec{x} \Theta^{-1} = \vec{x}, \qquad \Theta \vec{p} \Theta^{-1} = \vec{p}, \qquad \Theta \vec{J} \Theta^{-1} = -\vec{J}$$

$$| n \rangle = | n^{(0)} \rangle + | \delta n^{(1)} \rangle + | \delta n^{(2)} \rangle + \dots \qquad \Delta E_n^{(1)} = V_{nn}$$

$$| \delta n^{(1)} \rangle = \sum_{k \neq n} \frac{V_{kn}}{E_n^{(0)} - E_k^{(0)}} | k^{(0)} \rangle \qquad \Delta E_n^{(2)} = \sum_{k \neq n} \frac{|V_{nk}|^2}{E_n^{(0)} - E_k^{(0)}}$$

$$\dot{c}_n(t) = -\frac{i}{\hbar} \sum_m c_m(t) e^{i\omega_{nm}t} V_{nm} \qquad c_n^{(1)}(t) = -\frac{i}{\hbar} \int_0^t V_{ni} e^{i\omega_{ni}t'} dt'$$

$$c_n^{(2)}(t) = \left(-\frac{i}{\hbar} \right)^2 \sum_m \int_{t_0}^t dt' \int_{t_0}^{t'} dt'' V_{nm}(t') e^{i\omega_{nm}t'} V_{mi}(t'') e^{i\omega_{mi}t''}$$

$$| \psi^{(\pm)} \rangle = | \phi \rangle + \frac{1}{E - H_0 \pm i\epsilon} V | \psi^{(\pm)} \rangle \qquad f_\ell(k) = -\frac{\pi T_\ell(E)}{k}$$

$$f(\vec{k}', \vec{k}) = -\frac{1}{4\pi} \frac{2m}{\hbar^2} (2\pi)^3 \langle \vec{k}' | T | \vec{k} \rangle \qquad f(\vec{k}', \vec{k}) = \sum_{\ell=0}^{\infty} (2\ell+1) f_\ell(k) P_\ell(\cos \theta)$$

$$S_l(k) = e^{2i\delta_l} = 1 + 2ik f_l(k) \qquad \sigma_{\text{tot}} = \frac{4\pi}{k^2} \sum_l (2l+1) \sin^2 \delta_l$$

$$f^{(1)}(\vec{k}', \vec{k}) = -\frac{m}{2\pi\hbar^2} \int d^3x' e^{i\vec{q} \cdot \vec{x}'} V(\vec{x}') \qquad f^{(1)}(\theta) = -\frac{2m}{\hbar^2 q} \int dr r V(r) \sin(qr)$$

$$A_l(r) = e^{i\delta_l} [\cos \delta_l j_l(kr) - \sin \delta_l n_l(kr)]$$

$$j_0(x) = \sin(x)/x \qquad n_0(x) = -\cos(x)/x$$

$$j_l(x) \xrightarrow{x \rightarrow \infty} \sin(x - l\pi/2)/x \qquad n_l(x) \xrightarrow{x \rightarrow \infty} -\cos(x - l\pi/2)/x$$